Stochastic PDEs on Hilbert space with irregular noise coefficients With a first part on McKean-Vlasov equations

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Part 1. McKean-Vlasov equations

- Solving McKean-Vlasov equation with various interaction coefficients via relative entropy.
- Smoothness of density in the case of an interaction kernel.
- Uniform in time propagation of chaos with a sharp rate

McKean-Vlasov SDEs with non-Lipschitz interaction

Intensive research in recent years on solving McKean-Vlasov type equations given $b:[0,\infty)\times\mathbb{R}^d\to\mathbb{R}^d$,

$$dX_t = \langle \mu_t, b(t, X_t, \cdot) \rangle dt + dW_t, \quad \mu_t = \text{Law}(X_t)$$
(1)

given W a $d\mbox{-dimensional}$ Brownian motion or stable Lévy process.

When b is Lipschitz continuous in its last variable, one may solve directly via a Gronwall argument.

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- When b is Lipschitz continuous in its last variable, one may solve directly via a Gronwall argument.
- ▶ When *b* is merely bounded measurable, one may write the drift as $B(t, X_t, \mu_t)$ with *B* continuous in μ in total variation distance in the sense that for any t > 0, $x \in \mathbb{R}^d$,

$$|B(t, x, \mu) - B(t, x, \nu)| \le C \|\mu - \nu\|_{TV},$$
(2)

Let $\Phi(\mu)$ be solution to (1) for each $\mu \in \mathcal{P}([0,T];\mathbb{R}^d)$, then

$$H(\Phi(\mu) \mid \Phi(\nu) = \frac{1}{2} \mathbb{E}^{P_{\mu}} \left[\int_{0}^{t} |B(s, X_{s}, \mu) - B(s, X_{s}, \nu)|^{2} ds \right]$$
(3)

3/29

Solution-continued

The right hand side becomes

$$\leq \frac{C^2}{2} \int_0^t \|\mu_s - \nu_s\|_{TV}^2 ds \leq C \int_0^T H(\mu_s \mid \nu_s) ds.$$
 (4)

The first inequality follows from Lipschitz continuity of B in the measure component, and the second follows from Pinsker's inequality. Then the existence and uniqueness of a solution follows from a fixed point argument.

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- This method has appeared in several different papers in the 2010s. See for example [Lacker,2018].
- We can solve McKean-Vlasov equation in a much wider generality, where the assumption

$$|B(t, x, \mu) - B(t, x, \nu)| \le C \|\mu - \nu\|_{TV},$$
(5)

is no longer valid.

Examples of new results

Linear growth, path dependent coefficients:

 $|b(t, x, y)| \le K(1 + ||x||_t + ||y||_t), \quad t \in [0, T]$ (6)

we can solve (extending Lacker 2021)

$$dX_t = \langle \mu, b(t, X, \cdot) \rangle dt + dW_t, \quad \mu = \text{Law}(X)$$
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Singular and linear growth coefficients: given $\frac{d}{p_1} + \frac{2}{q_1} < 1$,

$$b_1(t, x, y) \le h_t(x - y)$$
 for some $h \in L_t^{q_1}([0, T], L_x^{p_1}(\mathbb{R}^d))$,

$$\sup_{t,y} |b_2(t,x,y)| \le K(1+|x|^\beta) \text{ for } K > 0, \beta \in [0,1),$$
 (8)

we can solve (extending Röckner-Zhang 2019)

$$dX_t = \langle \mu_t, b_1 + b_2(t, X_t, \cdot) \rangle dt + dW_t, \quad \mu_t = \text{Law}(X_t)$$
 (9)

 b_1 must be state dependent, but b_2 can be path dependent.

Examples for fractional Brownian driving noise

(Extending Galeati, Harang, Mayorcas 2021 and other works)

• Assume
$$H \in (0, \frac{1}{2})$$
 and $\alpha > 1 - \frac{1}{2H}$, and

 $\|B(t,\cdot,\mu) - B(t,\cdot,\nu)\|_{B^{\alpha}_{\infty,\infty}} \lesssim \|\mu - \nu\|_{TV}, \mu, \nu \in \mathcal{P}(\mathbb{R}^d),$ (10)

we can solve, via Girsanov transform for FBMs,

$$dX_t = B(t, X_t, \mu_t)dt + dB_t^H, \quad \text{Law}(X_t) = \mu_t, \qquad (11)$$

and when the interaction has linear growth and path dependent, i.e. $|b(t, x, y)| \le K(1 + ||x||_t + ||y||_t)$, solve

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• When $H \in (\frac{1}{2}, 1)$, and $\beta > H - \frac{1}{2} > 0$, assume

$$|b(t, x, x') - b(s, y, y')| \lesssim (|x - y|^{\alpha} + |x' - y'|^{\alpha} + |t - s|^{\beta}),$$

we can then solve the state dependent version of (12).

Examples for SPDEs

Stochastic heat equation on [0, 1], f bounded measurable,

$$dY(t,\sigma) = dW(t) + \kappa \frac{\partial^2}{\partial \sigma^2} Y(t,\sigma) dt + \int \mathcal{L}_{Y(t)}(dZ) \int_0^1 f(Y(t,\sigma), Z(\sigma')) d\sigma' dt.$$
(13)

and its more abstract version, assuming G is Lipschitz in μ_t in total variation,

$$\frac{\partial}{\partial t}Y(t) = \frac{\partial^2}{\partial \sigma^2}Y(t)dt + G(t, Y(t), \mu_t)dt + dW(t), \quad (14)$$

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For (13), we can recover the O(k²/n²) rate of propagation of chaos in relative entropy (Lacker 2021).

Questions

What if the diffusion coefficient also depends on the measure μ ?

$$dX_t = B(t, X_t, \mu_t)dt + \sigma(X_t, \mu_t)dW_t, \quad \text{Law}(X_t) = \mu_t$$
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- Huang, Ren and Wang, arXiv:2304.07562.

Smoothness of density for McKean-Vlasov SDEs Given $b : \mathbb{R}^d \to \mathbb{R}^d$ bounded measurable, the SDE

$$dX_t = b(X_t)dt + dW_t \tag{17}$$

has a density but is quite irregular, while the McKean-Vlasov SDE

$$dX_t = \langle \mu_t, b(X_t - \cdot) \rangle dt + dW_t, \quad \text{Law}(X_t) = \mu_t$$
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has a much smoother density. Possible ways to see this:

Malliavin calculus: when b is at least C¹, we can fix μ and show μ_t has some Besov regularity via Malliavin calculus. Then bootstrap to prove μ_t has a smooth density.

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- A more direct approach, better use of Besov space norms with exponent p, get smoothness of density for short time and very irregular b (Hao, Röckner and Zhang, arxiv 2302.04392).

Related smoothing phenomenon

This phenomenon has appeared in other situations such as

- 2D Navier-Stokes equation with (degenerate) additive white noise forcing, [Mattingly and Pardoux 2005].
- The mean field convolution structure is critical for solving

$$dX_t = \langle \mu_t, b(X_t - \cdot) \rangle dt + dW_t, \quad \text{Law}(X_t) = \mu_t$$
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for distributional b in the regularity class $b \in C_b^{-1+\epsilon}$, $\epsilon > 0$. (de Raynal, Jabir, Menozzi arXiv:2205.11866).

Propagation of chaos with a sharp rate

Let $P_t^{n,k}$ be the k-marginal density of a weakly interacting diffusion process with n components,

$$dX_t^{n,i} = \frac{1}{n-1} \sum_{j \neq i} b(X_t^{n,i}, X_t^{n,j}) dt + dW_t^i.$$
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Let also μ_t be the law of the (limiting) McKean-Vlasov equation. Then

It is classically understood that

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In [Jabin-Wang 18] they showed this convergence rate for the vorticity formulation of 2D Navier-Stokes equation on the torus, with sufficiently smooth initial condition.

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If b is Lipschitz continuous or bounded measurable, then [Lacker, 2021] showed that we indeed have a O(k/n) rate for k << n, which is sharp in several cases.</p>

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- An unsuccessful approach to get the sharp O(k/n) rate for singular interactions.
- May use modulated free energy instead of relative entropy. In the very recent work [De Courcel, Rosengweig, Serfaty. Arxiv: 2304.05315], they prove uniform-in-time mean-field convergence for singular periodic Riesz flows (s<d on T^d) in the gradient case with a sharp rate in the modulated energy pseudo distance.

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- Have to assume bounded or Lipschitz interactions.
- Could be relaxed if the flow has more structure (Riesz flow?) and we use modulated free energy instead.

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- We explore interesting applications of the coupling method to solution theory of SPDEs.

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- Strong solutions can be proved when the solution is real valued. "On quasi-linear stochastic partial differential equations", Gyöngy and Pardoux, 1993.

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- Strong uniqueness: imposing Sobolev regularity conditions on *σ* and use PDE theory.
- None of them works in infinite dimensions.

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- ► Evolution equations on Hilbert spaces: Da Prato-Zabczyk. Consider orthogonal basis $(e_n(x))_{n \in \mathbb{N}}$,

$$W(dxdt) = \sum_{n=1}^{\infty} e_n(x) dB_t^{(n)} dx.$$

$$N(t,x) = \sum_{n=1}^{\infty} \int_0^t \int_{\mathbb{R}^d} \mathcal{S}(t-s,x-y)g(s,y)e_n(y)dydB_s^{(n)}.$$

Notion of solutions to stochastic PDEs

- ► Random field solutions/ martingale measures approach: JB. Walsh. Regard s and y on an equal footing. W(A) = ∫_A W(dyds).
- ► Evolution equations on Hilbert spaces: Da Prato-Zabczyk. Consider orthogonal basis $(e_n(x))_{n \in \mathbb{N}}$,

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 Da Prato-Debussche technique, rough path theory, regularity structure, paracontrolled calculus.

$$\partial_t \Phi = \Delta \Phi + C \Phi - \Phi^3 + \xi$$

Only for additive noise or sufficiently regular noise coefficient.

Known results for the stochastic heat equation

• Additive noise case $\partial_t u = \Delta u + f(u) + \frac{\partial^2 W}{\partial_t \partial_x}$.

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• Multiplicative, Hölder continuous noise coefficient: σ being $\frac{3}{4} + \epsilon$ -Hölder continuous in X(t, x), the random field case:

$$\frac{\partial}{\partial t}X(t,x) = \frac{1}{2}\Delta X(t,x)dt + \sigma(t,x,X(t,x))dW(t,x)$$

"Pathwise Uniqueness for Stochastic Heat Equations with Hölder Continuous Coefficients: the White Noise Case", 77 p".

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- ▶ *B* has linear growth.

Well-Posedness

Theorem (Well-posedness of Stochastic Heat equation) Under the assumptions in the previous slide, there exists a unique (probabilistic weak) mild solution to

$$dX_t = AX_t dt + B(X_t)dt + \sigma(X_t)dW_t, \quad X_0 \in H.$$
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Given $\frac{1}{2} + \epsilon$ -Hölder $F : H \to H$, unique weak-mild solution to

 $dX_t = AX_t dt + (-A)^{1/2} F(X_t) dt + B(X_t) dt + \sigma(X_t) dW_t,$ (26)

• Examples: Burgers type equations, $\xi \in (0, 2\pi)$

$$du(t,\xi) = \frac{\partial^2}{\partial\xi^2}u(t,\xi)dt + \frac{\partial}{\partial\xi}h(u(t,\xi))dt + \sigma(u(t,\xi))dW_t(\xi).$$

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Cahn-Hilliard equations in dimensions 1,2,3:

$$du(t,\xi) = -\Delta_{\xi}^2 u(t,\xi) dt + \Delta_{\xi} h(u(t,\xi)) dt + \sigma(u(t,\xi)) dW_t(\xi)$$

Long-time behaviour

Theorem (Exponential ergodicity)

Assume the drift $B : H \to H$ is Hölder continuous, and the Lyapunov condition hold: for some $V : H \to \mathbb{R}_+$ and some $\lambda \in (0, 1)$ infinity at infinity,

$$\mathbb{E}[V(X_t)] \le \lambda V(X_0) + M \tag{27}$$

for some given t > 0 and M > 0. Then there exists a unique invariant measure, and the solution converges to the invariant measure exponentially fast with respect to (some specific) Wasserstein distance on $\mathcal{P}(H)$.

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- No applicable Itô formula for cylindrical noise
- Lyapunov assumption satisfied when B, F, σ are bounded, and A is a negative operator.

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• Use pathwise estimate to compare $X|_{[0,T]}$ with $\widetilde{X}^n|_{[0,T]}$.
Some technical challenges in infinite dimensions:

• Derive maximal inequality for the process, when λ is large:

 $dX_t = \Delta X_t dt - \lambda X_t dt + \Phi(t) dW_t, \quad X_0 = 0,$

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- No applicable Itô's formula: always work with mild formulations.
- Lipschitz approximation in infinite dimensions: compactness of heat semigroup.

Now I outline the procedure of proof:

• Compactness reduction: for any $\epsilon > 0$ find a compact subset $K \subset H$ s.t. X_t stays in K for $t \in [0,T]$, w.p. at least $1 - \epsilon$

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- ► Combine pathwise estimates between X_t and X̃ⁿ_t, and probabilistic estimates between Xⁿ_t and X̃ⁿ_t. This gives a measurement of how close X_t is to Xⁿ_t.

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- ▶ Use also the estimate $\|S(t)A^{1/2}\|_{op} \leq C\frac{1}{\sqrt{t}}$ in the presence of the $(-A)^{1/2}$ term in the drift.

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- choose $\lambda = c_n^{\gamma-1}$ for some $\gamma > 0$, and choose the stopping time $\tau \in [0, T]$ the first time that $|X_t \widetilde{X}_t^n| \ge c_n$.
- ► Combine pathwise estimates between X_t and X_tⁿ, and probabilistic estimates between X_tⁿ and X_tⁿ. This gives a measurement of how close X_t is to X_tⁿ.
- Use also the estimate $\|S(t)A^{1/2}\|_{op} \leq C\frac{1}{\sqrt{t}}$ in the presence of the $(-A)^{1/2}$ term in the drift.
- Proof of ergodic behavior follows similar lines but need an extra argument constructing the Wasserstein distance on H.

Stochastic wave equation: well-posedness

Our method works not only for the parabolic systems, but also for hyperbolic systems.

Consider the (abstract) damped stochastic wave equation

$$\mu \frac{\partial^2 u_{\mu}(t)}{\partial t^2} = A u_{\mu}(t) - \frac{\partial u_{\mu}(t)}{\partial t} + B(t, u_{\mu}(t)) + G(t, u_{\mu}(t)) dW_t,$$
(28)

and the stochastic wave equation without damping term

$$\mu \frac{\partial^2 u_{\mu}(t)}{\partial t^2} = A u_{\mu}(t) + B(t, u_{\mu}(t)) + G(t, u_{\mu}(t)) dW_t,$$
(29)

Theorem (Well-posedness of stochastic wave equation) Under the same assumption on A, B and G as in the case of the stochastic heat equation, there exists a unique weak-mild solution to (28) and (29).

Stochastic wave equation: small mass limit

Theorem

Assume moreover that B is Hölder continuous in u_{μ} . Then as μ tends to 0, the solution to the damped stochastic wave equation

$$\mu \frac{\partial^2 u_{\mu}(t)}{\partial t^2} = A u_{\mu}(t) - \frac{\partial u_{\mu}(t)}{\partial t} + B(t, u_{\mu}(t)) + G(t, u_{\mu}(t)) dW_t$$

converges in distribution on path space to the solution of the stochastic heat equation

$$\frac{\partial u(t)}{\partial t} = Au(t) + B(t, u(t)) + G(t, u(t))dW_t.$$

"On the Smoluchowski-Kramers approximation for a system with an infinite number of degrees of freedom", Freidlin and Cerrai, 2006.

Some remaining questions to be addressed:

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- In the presence of (−A)^{1/2}F term, we require F to be ¹/₂ + ε-Hölder. Can we allow for ε-Hölder? [By Priola 2021, we can have ε-Hölder continuity in the case of additive noise.]
- Need a better understanding of infinite dimensional Kolmogorov equation beyond the well studied case of additive noise.

We discussed the general setting of Hilbert space valued solutions, but some better estimates hold for random field solutions.

Consider the parabolic stochastic PDE

$$\partial_t u(t,x) = \partial_x^2 u(t,x) + g(t,x,u) + \sigma(t,x,u) dW(t,x)$$

where $x \in [0,1], \, W$ is the space-time white noise on [0,1] and g is uniformly bounded.

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Thanks